Yangian symmetry of fishnet graphs

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Overview

• Planar multi-loop multi-point massless conformal Feynman graphs



[Zamolodchikov '80]

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- Yangian symmetry \longleftrightarrow Integrability
- Higher order symmetry (extends conformal symmetry)
- Set of differential equations
- Integrability of the underlying theory
- 4D biscalar theory (limit of $\mathcal{N} = 4$ SYM) \iff square fishnet lattice
 - Single-trace correlators (off-shell legs)
 - Amplitudes (on-shell legs)
 Yangian symmetry
 - Cuts (mixed on-shell/off-shell)
- Yangian symmetry is NOT broken by loop corrections
- Generalization: 3D, 6D scalar theory, 4D scalars & fermions

- QFT generating Fishnet graphs (square lattice)
- 4D biscalar theory. Complex scalars ϕ_1, ϕ_2 in adj of $SU(N_c)$

$$\mathcal{L} = \frac{N_c}{2} \operatorname{Tr} \left(\partial^{\mu} \phi_1^{\dagger} \partial_{\mu} \phi_1 + \partial^{\mu} \phi_2^{\dagger} \partial_{\mu} \phi_2 + 2\xi^2 \phi_1^{\dagger} \phi_2^{\dagger} \phi_1 \phi_2 \right)$$

[Gurdogan, Kazakov '15]

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- A double scaling limit of γ -deformed $\mathcal{N}=4$ SYM
- "Almost" conformal in the planar limit. $\frac{d\xi}{d\log \mu} = O(N_c^{-2})$
- Integrability (spectrum of anomalous dimensions) [Caetano, Gurdogan, Kazakov '16]
- Non-unitary. Chiral structure

$$\phi_2 \longrightarrow \phi_1^{\dagger} \phi_2^{\dagger}$$

One Feynman graph per loop order in the planar limit

Fishnet graphs with regular boundary



 $\langle \operatorname{Tr} [\chi_1(x_1) \chi_2(x_2), \ldots \chi_{2M}(x_{2M})] \rangle$

where $\chi_i \in \{\phi_1^{\dagger}, \phi_2^{\dagger}, \phi_1, \phi_2\}$





Duality transformation

$$p_i^{\mu} = x_i^{\mu} - x_{i+1}^{\mu}$$

x-space – correlator graph

p-space – loop graph with off-shell legs

Yangian of conformal algebra

Conformal algebra $\mathfrak{so}(2,4)$

$$D = -i(x_{\mu}\partial_{\mu} + \Delta) \quad , \quad L_{\mu\nu} = i(x_{\mu}\partial_{\nu} - x_{\nu}\partial_{\mu}) \; ,$$
$$P_{\mu} = -i\partial_{\mu} \quad , \quad K_{\mu} = i(x^{2}\partial_{\mu} - 2x_{\mu}x^{\nu}\partial_{\nu} - 2\Delta x_{\mu})$$

and its infinite-dimensional extension – Yangian

$$J \ , \ \widehat{J} \ , \ [\widehat{J}, \widehat{J}] \ , \ [\widehat{J}, \widehat{J}] \ , \ [\widehat{J}, [\widehat{J}, \widehat{J}]] \ , \ [[\widehat{J}, \widehat{J}], [\widehat{J}, \widehat{J}]] \ \ldots$$

where level-zero generators $J^A \in \mathfrak{so}(2,4)$ and level-one generators \widehat{J} satisfy

$$\begin{bmatrix} J^A, J^B \end{bmatrix} = f^{AB}{}_C J^C \quad , \quad \begin{bmatrix} J^A, \hat{J}^B \end{bmatrix} = f^{AB}{}_C \hat{J}^C \quad , \quad \begin{array}{c} \text{Jacobi and Serre} \\ \text{relations} - \text{cubic in } J, \hat{J} \end{bmatrix}$$

Evaluation representation with evaluation parameters v_k

$$\widehat{J}^{A} = \frac{1}{2} f^{A}_{BC} \sum_{j < k} J^{C}_{j} J^{B}_{k} + \sum_{k} v_{k} J^{A}_{k}$$

Yangian symmetry of the Fishnet graphs

$$J^{\mathcal{A}}|\mathsf{Fishnet}
angle=\widehat{J}^{\mathcal{A}}|\mathsf{Fishnet}
angle=0$$

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[Drinfeld '85]

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Evaluation representation with evaluation parameters v_k

$$\widehat{P}^{\mu} = -\frac{i}{2} \sum_{j < k} \left[(\mathcal{L}_{j}^{\mu\nu} + \eta^{\mu\nu} \mathcal{D}_{j}) \mathcal{P}_{k,\nu} - (j \leftrightarrow k) \right] + \sum_{k} \mathsf{v}_{k} \mathcal{P}_{k}^{\mu}$$

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[Drinfeld '85]

Lax and monodomry matrix

• Lax matrix with the spectral parameter u

$$L(u) = \begin{pmatrix} 1 + \frac{1}{u}\mathcal{J}_{11} & \frac{1}{u}\mathcal{J}_{12} & \frac{1}{u}\mathcal{J}_{13} & \frac{1}{u}\mathcal{J}_{14} \\ \frac{1}{u}\mathcal{J}_{21} & 1 + \frac{1}{u}\mathcal{J}_{22} & \frac{1}{u}\mathcal{J}_{23} & \frac{1}{u}\mathcal{J}_{24} \\ \frac{1}{u}\mathcal{J}_{31} & \frac{1}{u}\mathcal{J}_{32} & 1 + \frac{1}{u}\mathcal{J}_{33} & \frac{1}{u}\mathcal{J}_{34} \\ \frac{1}{u}\mathcal{J}_{41} & \frac{1}{u}\mathcal{J}_{42} & \frac{1}{u}\mathcal{J}_{43} & 1 + \frac{1}{u}\mathcal{J}_{44} \end{pmatrix}$$

consists of $\mathfrak{so}(2,4)$ generators $\mathcal{J}_{ij} \in \operatorname{span}\{D, P_{\mu}, K_{\nu}, L_{\mu\nu}\}$

• n-point monodromy matrix with inhomogeneities $\delta_1, \delta_2, \ldots, \delta_n$

$$T(u; \vec{\delta}) = L_n(u + \delta_n) \dots L_2(u + \delta_2) L_1(u + \delta_1)$$

$$T_{ab}(u; \vec{\delta}) = \delta_{ab} + \sum_{k \ge 0} u^{-1-k} \mathcal{J}_{ab}^{(k)}$$

Quantum spin chain with noncompact representations of $\mathfrak{so}(2,4)$

Lax and monodomry matrix

RTT-relation defines the quadratic algebra for $\{\mathcal{J}^{(k)}\}$

$$R_{ae,bf}(u-v) T_{ec}(u) T_{fd}(v) = T_{ae}(v) T_{bf}(u) R_{ec,fd}(u-v)$$

[Faddeev, Kulish, Sklyanin, Takhtajan,...'79]

with Yang's R-matrix $R_{ab,cd}(u) = \delta_{ab}\delta_{cd} + u\delta_{ad}\delta_{bc}$. RTT is compatible with the co-product.

Yangian symmetry \leftrightarrow Eigenvalue relation for the monodromy matrix

$$L_n(u+\delta_n)\dots L_2(u+\delta_2)\;L_1(u+\delta_1)\;|\mathsf{Fishnet}
angle=\lambda(u;ec\delta)\,|\mathsf{Fishnet}
angle\cdot\mathbf{1}$$

and expanding this matrix equation in the spectral parameter u,

$$|\mathcal{J}_{ab}^{(k)}|\mathsf{Fishnet}
angle=\lambda_{n}(ec{\delta})\delta_{ab}\,|\mathsf{Fishnet}
angle$$

For oscillator representations [D.C., Kirschner '13; D.C., Kirschner, Derkachov '13; Frassek, Kanning, Ko, Staudacher '13; Broedel, de Leeuw, Rosso '14; ...]

Conformal Lax

Lax matrix depends on parameters $(u,\Delta) \Leftrightarrow (u_+,u_-)$

$$L(u_{+}, u_{-}) = \begin{pmatrix} u_{+} \cdot \mathbf{1} - \mathbf{p} \cdot \mathbf{x} & \mathbf{p} \\ -\mathbf{x} \cdot \mathbf{p} \cdot \mathbf{x} + (u_{+} - u_{-}) \cdot \mathbf{x} & u_{-} \cdot \mathbf{1} + \mathbf{x} \cdot \mathbf{p} \end{pmatrix}$$

where

[D.C., Derkachov, Isaev '12]

$$\mathbf{x} = -i\overline{\sigma}_{\mu}\mathbf{x}^{\mu}$$
, $\mathbf{p} = -\frac{i}{2}\sigma^{\mu}\partial_{\mu}$, $u_{+} = u + \frac{\Delta-4}{2}$, $u_{-} = u - \frac{\Delta}{2}$

• Local vacuum of the Lax

$$L(u, u+2) \cdot 1 = (u+2) \cdot \mathbf{1}$$

• Intertwining relation with the scalar propagator

 $x_{i,j} \equiv x_i - x_j$

 $\sum_{x_{1}}^{x_{2}} \sum_{L_{2}(\bullet, u+1)}^{L_{2}(\bullet, u+1)} = \int_{L_{1}(u, *)}^{x_{2}} \sum_{L_{1}(u, *)}^{L_{2}(\bullet, u+1)} \sum_{L_{1}(u, *)}^{L_{1}(u+1, *)L_{2}(\bullet, u)} x_{12}^{-2}$ $= x_{12}^{-2} L_{1}(u, *)L_{2}(\bullet, u+1)$ • Integration by parts $L^{T}(u+2, u) \cdot 1 = (u+2) \cdot 1$

Example: Cross integral









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 $L^{T}(u+2,u)\cdot 1 = (u+2)\cdot \mathbf{1} \qquad ,$

Example: Cross integral



 $L^{T}(u+2, u) \cdot 1 = (u+2) \cdot \mathbf{1} , \qquad L(u, u+2) \cdot 1 = (u+2) \cdot \mathbf{1}$ $x_{12}^{-2}L_{1}(u, *)L_{2}(\bullet, u+1) = L_{1}(u+1, *)L_{2}(\bullet, u)x_{12}^{-2}$

Example: Cross integral



 $L^{T}(u+2, u) \cdot 1 = (u+2) \cdot \mathbf{1} , \qquad L(u, u+2) \cdot 1 = (u+2) \cdot \mathbf{1}$ $x_{12}^{-2}L_{1}(u, *)L_{2}(\bullet, u+1) = L_{1}(u+1, *)L_{2}(\bullet, u)x_{12}^{-2}$

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Example: Cross integral



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Example: Cross integral

Yangian symmetry of the cross integral

 $L_{4}[4,5]L_{3}[3,4]L_{2}[2,3]L_{1}[1,2]|cross\rangle = [3][4]^{2}[5] \cdot |cross\rangle \cdot \mathbf{1}$ where $[i,j] \equiv L(u+i,u+j)$ and $[i] \equiv u+i$ $|cross\rangle = \int d^{4}x_{0}\frac{1}{x_{10}^{2}x_{20}^{2}x_{30}^{2}x_{40}^{2}} = x_{13}^{-2}x_{24}^{-2}\Phi(s,t)$

Conformal cross-ratios s, t. The Yangian symmetry implies DE

$$\Phi + (3s-1)\frac{\partial \Phi}{\partial s} + 3t\frac{\partial \Phi}{\partial t} + (s-1)s\frac{\partial^2 \Phi}{\partial s^2} + t^2\frac{\partial^2 \Phi}{\partial t^2} + 2st\frac{\partial^2 \Phi}{\partial s\partial t} = 0$$

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Fishnet graphs with regular boundary



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where $\mathcal{C} = \mathcal{C}_{\text{in}} \cup \mathcal{C}_{\text{out}}$ is the boundary of the graph

Fishnet graphs with regular boundary



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where $\mathcal{C} = \mathcal{C}_{\text{in}} \cup \mathcal{C}_{\text{out}}$ is the boundary of the graph

Fishnet graphs with irregular boundary



Single-trace correlator

$$\langle \operatorname{Tr} [\chi_1(x_1) \chi_2(x_2), \ldots \chi_{2M}(x_{2M})] \rangle$$

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where $\chi_i \in \{\phi_1^{\dagger}, \phi_2^{\dagger}, \phi_1, \phi_2\}$ and some of x_i 's are identified

Fishnet graphs with irregular boundary



and eigenvalue λ is fixed by algebraic equations (implied by cyclic symmetry).

Scattering amplitudes and cuts of fishnet graphs

Duality transformation $x_i^{\mu} - x_{i+1}^{\mu} = p_i^{\mu}$. Amplitude $p_i^2 = 0 \Leftrightarrow x_{i,i+1}^2 = 0$



 $T(u) |\mathsf{Fishnet}\rangle = \lambda(u) |\mathsf{Fishnet}\rangle \cdot \mathbf{1} \quad \Rightarrow \quad T(u) |\mathsf{Fishnet}\rangle_{\mathsf{cut}} = \lambda(u) |\mathsf{Fishnet}\rangle_{\mathsf{cut}} \cdot \mathbf{1}$ for any cut $x_{ij}^{-2} \rightarrow \delta(x_{ij}^2)$.

Generalizations: tri-scalar theories in 3D, 6D, and scalar-fermion theories in 4D

$$\mathcal{L}_{3D,\mathrm{int}} = \xi N_c \mathrm{Tr} Y_1 Y_4^{\dagger} Y_2 Y_1^{\dagger} Y_4 Y_2^{\dagger}$$

[Caetano, Gurdogan, Kazakov '16]



$$\mathcal{L}_{4D,\mathrm{int}} = \mathcal{L}_{\mathsf{Yukawa}} + \mathcal{L}_{\phi^4}$$

scalar and fermion fields

[Gurdogan, Kazakov '15]



$$\mathcal{L}_{6D,\mathrm{int}} = N_c \mathrm{Tr} \left(\xi_1 \, \phi_1^{\dagger} \phi_2 \phi_3 + \xi_2 \, \phi_1 \phi_2^{\dagger} \phi_3^{\dagger} \right)$$

[Mamroud, Torrents '17]

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Conclusions

- Conformal Yangian symmetry of single-trace correlators and amplitudes
- RTT-realization of the Yangian algebra
- Propagators intertwining operators of the Laxes and monodromy
- Yangian invariance is a consequence of (bi)local operator relations
- Lasso proof of the Yangian invariance
- $\bullet\,$ The approach also works for Fishnet graphs in 3D, 6D, and for scalar-fermion theories in 4D